

ASYMPTOTIC CURVATURE DECAY AND REMOVAL OF SINGULARITIES OF BACH-FLAT METRICS

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ABSTRACT. We prove a removal of singularities result for Bach-flat metrics in dimension 4 under the assumption of bounded L^2 norm of curvature, bounded Sobolev constant and a volume growth bound. This result extends the removal of singularities result for special classes of Bach-flat metrics obtained in [6]. For the proof we analyze the decay rates of solutions to the Bach-flat equation linearized around a flat metric. This classification is used to prove that Bach-flat cones are in fact ALE of order τ for any $\tau < 2$. This result is then used to prove the removal of singularities theorem.

1. INTRODUCTION

In dimension 4, the Euler-Lagrange equations for the functional

$$(1.1) \quad \mathcal{W}(g) := \int_M |W_g|^2 dV_g$$

where W_g is the Weyl curvature, are given by

$$(1.2) \quad B_{ij} = \nabla^k \nabla^l W_{ikjl} + \frac{1}{2} R^{kl} W_{ikjl} = 0$$

where W_{ijkl} and R^{kl} correspond to the components of the Weyl and Ricci tensors respectively [2], [4]. B_{ij} are the components of the *Bach tensor*, and a metric which is critical for \mathcal{W} is called *Bach-flat*.

A smooth Riemannian manifold (X^4, g) is called an asymptotically locally Euclidean end of order τ if there exists a finite subgroup $\Gamma \subset SO(4)$ which acts freely on $\mathbb{R}^4 \setminus B(0, R)$ and a C^∞ diffeomorphism $\phi : X \rightarrow (\mathbb{R}^4 \setminus B(0, R)) / \Gamma$ such that using this identification

$$(1.3) \quad g_{ij} = \delta_{ij} + O(r^{-\tau})$$

$$(1.4) \quad \partial^{[k} g_{ij} = O(r^{-\tau-k})$$

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for any partial derivative of order k as $r \rightarrow \infty$. We say an end is ALE of order 0 if there exist coordinates as above so that

$$(1.5) \quad g_{ij} = \delta_{ij} + o(1)$$

$$(1.6) \quad \partial^{|k|} g_{ij} = o(r^{-k})$$

as $r \rightarrow \infty$.

Theorem: 1.1. *Let (X^4, g) be a complete, noncompact four-dimensional Bach-flat space with zero scalar curvature which is ALE of order 0. Assume further that*

$$(1.7) \quad \int_X |\text{Rm}_g|^2 dV_g < \Lambda < \infty, C_S < K < \infty$$

Then the cone is ALE of order τ for any $\tau < 2$.

This theorem is an extension of theorem 1.1 of [7], where under the same hypotheses on curvature and Sobolev constant and also assuming $b_1(X) < \infty$ where $b_1(X)$ is the first Betti number it is proved that the end is ALE of order zero.

Theorem: 1.2. *Let (X, d, x) be a complete locally compact length space with base point x . Let $B(x, 1) \setminus \{x\}$ be a connected C^∞ Bach-flat four-manifold satisfying*

$$(1.8) \quad \int_{B(x, 1)} |\text{Rm}_g|^2 dV_g < \infty$$

$$(1.9) \quad \|u\|_{L^4(B(x, 1) \setminus \{x\})} \leq C_s \|\nabla u\|_{L^2(B(x, 1) \setminus \{x\})}, u \in C^{0,1}(B(x, 1) \setminus \{x\})$$

$$(1.10) \quad \text{Vol}(B(x, r)) \leq V_1 r^4, r > 0$$

$$(1.11) \quad b_1(X) < \infty$$

where C_s and V_1 are positive constants. Then the metric g extends to $B(x, 1)$ as a smooth orbifold metric.

This is an extension of theorem 6.4 of [6], and indeed this result was suggested in the work of Gang Tian and Jeff Viaclovsky [6], [7]. There removal of singularities is proved for certain subclasses of Bach-flat manifolds, specifically half-conformally flat metrics, metrics with harmonic curvature, and constant scalar curvature Kähler metrics. The proof uses certain Kato inequalities which are satisfied in these special cases to improve the decay of the Ricci tensor, which in turn can be used to get improved decay of the full curvature tensor. The techniques

used herin also suffice to give this result assuming that the metric has harmonic curvature, that is,

$$(1.12) \quad \delta \text{Rm} = -\nabla^i R_{ijkl} = 0$$

We now give an outline of the rest of the paper. In section 2 we decompose the action of the Bilaplacian acting on symmetric two-tensors according to the radial separation of variables on a cone. In section 3 we compute the linearization of the Bach-flat condition around a flat metric. We observe that after applying a change of diffeomorphism gauge and possibly a conformal change this equation is equivalent to the equation $\Delta^2 h = 0$ where h is the perturbation of the metric. Using the separation of variables we further simplify this equation and classify the solutions in terms of their decay or growth rates at infinity. Section 4 gives the proof of theorem 1.1, which consists of considering the Bach-flat equation at infinity in the cone as a perturbation of a flat metric, and using our analysis of the linearized equation. Section 5 uses theorem 1.1 to prove theorem 1.2.

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2. DECOMPOSITION OF THE BILAPLACIAN

In this section we will decompose the action of the Bilaplacian on symmetric two-tensors on \mathbb{R}^4 according to the cone structure. These tensors naturally decompose according to radial and tangential directions. We reintroduce the notation of [1] and recall the decomposition of the Laplacian proved there. These lemmas are then used in a straightforward way to decompose the action of the Bilaplacian. The calculation is greatly simplified by restricting to \mathbb{R}^4 , where we completely understand the curvature tensor, and how it decomposes according to the cone structure.

Write $(\mathbb{R}^4, g_{\text{flat}})$ as $C(S^3, g_{\text{can}})$, the cone over the three-sphere with its canonical metric. Let $P : T(1, S^3) \rightarrow T(r, S^3)$ denote the identification of tangent bundles induced by parallel transport along radial geodesics, and let $\tilde{\nabla}^r$ denote the Riemannian connection with respect to the induced metric on (r, S^3) . Define the connection

$$(2.1) \quad \tilde{\nabla} = P \tilde{\nabla}^1 P^{-1}$$

A basic formula is then

$$(2.2) \quad \tilde{\nabla}^r = r^{-1} \tilde{\nabla}$$

Also let $\tilde{d}^r, \tilde{\delta}^r, \tilde{\text{tr}}^r$ be defined with respect to the induced metric on (r, S^3) and let \tilde{d}^r denote exterior differentiation on (r, S^3) . Then define

$$\begin{aligned} \tilde{\underline{\delta}} &= P\tilde{\delta}^1 P^{-1} \\ \tilde{\underline{d}} &= P\tilde{d}^1 P^{-1} \\ \tilde{\underline{\text{tr}}} &= P\tilde{\text{tr}}^1 P^{-1} \end{aligned} \quad (2.3)$$

The following formulae are immediate

$$\begin{aligned} \tilde{\Delta}^r &= r^{-2} \tilde{\underline{\Delta}} \\ \tilde{\delta}^r &= r^{-1} \tilde{\underline{\delta}} \\ \tilde{d}^r &= r^{-1} \tilde{\underline{d}} \\ \tilde{\text{tr}}^r &= \tilde{\underline{\text{tr}}} \end{aligned} \quad (2.4)$$

Now say e is a vector tangent to (r, S^3) and e^* is its dual 1-form. Let η be a 1-form such that

$$\eta \left(\frac{\partial}{\partial r} \right) = 0 \quad (2.5)$$

$$\nabla_{\frac{\partial}{\partial r}} \eta \equiv 0 \quad (2.6)$$

Some simple formulas are then

$$\nabla_e \eta = r^{-1} \left(\tilde{\underline{\nabla}}_e \eta - \eta(e) dr \right) \quad (2.7)$$

$$\nabla_e dr = r^{-1} e^* \quad (2.8)$$

$$\nabla_{\frac{\partial}{\partial r}} \frac{\partial}{\partial r} = \nabla_{\frac{\partial}{\partial r}} dr = 0 \quad (2.9)$$

Now fix $\frac{\partial}{\partial r}, e_1, e_2, e_3$ a local orthonormal basis satisfying

$$\tilde{\nabla}_{e_i}^r e_j = 0 \quad (2.10)$$

at a fixed point (r, x) and also

$$\nabla_{\frac{\partial}{\partial r}} e \equiv 0 \quad (2.11)$$

Finally, given ω_1, ω_2 1-forms let

$$\omega_1 \boxtimes \omega_2 = \omega_1 \otimes \omega_2 + \omega_2 \otimes \omega_1 \quad (2.12)$$

Using this frame the symmetric two-tensors naturally decompose into three distinct types. In the next four lemmas we record the action of divergence and the Laplacian on these tensors. We will need these preliminary calculations to calculate the action of the Bilaplacian on each of these types of tensors. The proofs of these lemmas can all be found in [1].

Lemma: 2.1. *Consider $f(r)B$ where $B = \sum \eta_i \boxtimes \eta_j$ and η_i satisfies (2.5) and (2.6). Then*

$$(2.13) \quad \begin{aligned} \nabla^* \nabla(fB) = & \left(-f'' - 3r^{-1}f' + r^{-2}f \left(\widetilde{\nabla}^* \widetilde{\nabla} + 2 \right) \right) B \\ & + 2r^{-2}f\widetilde{\underline{\delta}}(B) \boxtimes dr - 2r^{-2}f\widetilde{\underline{\text{tr}}}(B)dr \otimes dr \end{aligned}$$

Lemma: 2.2. *Consider $k(r)\tau \boxtimes dr$ where τ satisfies (2.5) and (2.6). Then*

$$(2.14) \quad \begin{aligned} \nabla^* \nabla(k\tau \boxtimes dr) = & -2r^{-2}k\widetilde{\nabla}^{\text{sym}}\tau \\ & + \left\{ \left[-k'' - 3r^{-1}k' + r^{-2}k \left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \right] \tau \right\} \boxtimes dr \\ & + 4r^{-2}k\widetilde{\underline{\delta}}\tau dr \otimes dr \end{aligned}$$

where

$$(2.15) \quad \widetilde{\nabla}^{\text{sym}}\tau = \sum_i \widetilde{\nabla}_{e_i}\tau \boxtimes e_i^*$$

Lemma: 2.3. *Consider $l(r)\phi(x)dr \otimes dr$. Then*

$$(2.16) \quad \begin{aligned} \nabla^* \nabla(l\phi dr \otimes dr) = & -2r^{-2}l\phi\widetilde{\underline{g}} - 2r^{-2}l\widetilde{\underline{d}}\phi \boxtimes dr \\ & + (-l'' - 3r^{-1}l' + r^{-2}l(\widetilde{\nabla}^* \widetilde{\nabla} + 6))\phi dr \otimes dr \end{aligned}$$

Lemma: 2.4. *For tensors of the three types considered above, we have the following formulas:*

$$(2.17) \quad \delta(fB) = r^{-1}f\widetilde{\underline{\delta}}(B) - r^{-1}f\widetilde{\underline{\text{tr}}}(B)dr$$

$$(2.18) \quad \delta(k\tau \boxtimes dr) = k'\tau + 4r^{-1}k\tau + r^{-1}k \left(\widetilde{\underline{\delta}}\tau \right) dr$$

$$(2.19) \quad \delta(l\phi dr \otimes dr) = (l' + 3r^{-1}l)\phi dr$$

Proposition: 2.5. *Consider $f(r)B$ where $B = \sum \eta_i \boxtimes \eta_j$ and η_i satisfies (2.5) and (2.6). Then*

$$\begin{aligned}
 (2.20) \quad (\nabla^* \nabla)^2 (fB) = & \left(f^{(4)} + 6r^{-1}f^{(3)} - r^{-2}f'' \left(1 + 2\tilde{\nabla}^* \tilde{\nabla} \right) \right) B \\
 & - \left(r^{-3}f' \left(2\tilde{\nabla}^* \tilde{\nabla} + 7 \right) \right) B + r^{-4}f \left(\left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) \right)^2 B \\
 & - 4r^{-4}f\tilde{\nabla}^{\text{sym}}\tilde{\delta}(B) + 4r^{-4}f\tilde{\text{tr}}(B)\tilde{g} \\
 & + 4 \left\{ \left[-r^{-2}f'' - r^{-3}f' + r^{-4}f \left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) \right] \tilde{\delta}(B) \right\} \boxtimes dr \\
 & + 8r^{-4}f\tilde{d}\tilde{\text{tr}}(B) \boxtimes dr \\
 & + 4 \left(r^{-2}f'' + r^{-3}f' - r^{-4}f \left(\tilde{\nabla}^* \tilde{\nabla} + 4 \right) \right) \tilde{\text{tr}}(B)dr \otimes dr \\
 & + 8r^{-4}f\tilde{\delta}\tilde{\delta}(B)dr \otimes dr
 \end{aligned}$$

Proof. This will be a straightforward calculation using the three above lemmas. First of all using (2.13) we see that

$$\begin{aligned}
 (2.21) \quad -\nabla^* \nabla (f''B) = & \left(f^{(4)} + 3r^{-1}f^{(3)} - r^{-2}f'' \left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) \right) B \\
 & - 2r^{-2}f''\tilde{\delta}(B) \boxtimes dr + 2r^{-2}f''\tilde{\text{tr}}(B)dr \otimes dr
 \end{aligned}$$

Next

$$\begin{aligned}
 (2.22) \quad \nabla^* \nabla (-3r^{-1}f'B) = & 3 \left(r^{-1}f^{(3)} + r^{-2}f'' - r^{-3}f' \left(\tilde{\nabla}^* \tilde{\nabla} + 3 \right) \right) B \\
 & - 6r^{-3}f'\tilde{\delta}(B) \boxtimes dr + 6r^{-3}f'\tilde{\text{tr}}(B)dr \otimes dr
 \end{aligned}$$

Now it is clear that $\tilde{\nabla}^* \tilde{\nabla} B$ satisfies (2.5) and (2.6). Thus apply (2.13) again to compute

$$\begin{aligned}
 (2.23) \quad \nabla^* \nabla \left(r^{-2}f \left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) B \right) = & \left(-r^{-2}f'' + r^{-3}f' + r^{-4}f \left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) \right) \left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) (B) \\
 & + 2r^{-4}f \left(\tilde{\nabla}^* \tilde{\nabla} - 2 \right) (\tilde{\delta}B) \boxtimes dr + 4r^{-4}f\tilde{d}\tilde{\text{tr}}B \boxtimes dr \\
 & - 2r^{-4}f \left(\tilde{\nabla}^* \tilde{\nabla} + 2 \right) \tilde{\text{tr}}(B)dr \otimes dr
 \end{aligned}$$

Note that we have applied lemma 6.2 to commute $\tilde{\delta}$ past $\tilde{\nabla}^* \tilde{\nabla}$. Now, since η satisfies (2.5) and (2.6), so does $\tilde{\delta}B$, thus we may apply (2.14)

to conclude

$$\begin{aligned}
 (2.24) \quad & \nabla^* \nabla \left(2r^{-2} f \tilde{\underline{\delta}}(B) \boxtimes dr \right) \\
 &= -4r^{-4} f \tilde{\underline{\nabla}}^{\text{sym}} \tilde{\underline{\delta}}(B) \\
 &\quad + \left\{ \left[-2r^{-2} f'' + 2r^{-3} f' + 2r^{-4} f \left(\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 6 \right) \right] \tilde{\underline{\delta}}(B) \right\} \boxtimes dr \\
 &\quad + 8r^{-4} f \tilde{\underline{\delta}} \tilde{\underline{\delta}}(B) dr \otimes dr
 \end{aligned}$$

Similarly we apply (2.16) to conclude

$$\begin{aligned}
 (2.25) \quad & \nabla^* \nabla \left(-2r^{-2} f \tilde{\underline{\text{tr}}}(B) dr \otimes dr \right) \\
 &= 4r^{-4} f \tilde{\underline{\text{tr}}}(B) \tilde{\underline{g}} + 4r^{-4} f \tilde{\underline{d}} \tilde{\underline{\text{tr}}}(B) \boxtimes dr \\
 &\quad + \left(2r^{-2} f'' - 2r^{-3} f' - 2r^{-4} f \left(\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 6 \right) \right) \tilde{\underline{\text{tr}}}(B) dr \otimes dr
 \end{aligned}$$

Summing together (2.21) - (2.25) gives the result. \square

Proposition: 2.6. *Consider $k(r)\tau \boxtimes dr$ where τ satisfies (2.5) and (2.6). Then*

$$\begin{aligned}
 (2.26) \quad & (\nabla^* \nabla)^2 (k\tau \boxtimes dr) \\
 &= 4 \left(r^{-2} k'' + r^{-3} k' - r^{-4} k \left(\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 5 \right) \right) \left(\tilde{\underline{\nabla}}^{\text{sym}} \tau \right) \\
 &\quad - 8r^{-4} k \tilde{\underline{\delta}} \tau \tilde{\underline{g}} \\
 &\quad + \left(k^{(4)} + 6r^{-1} k^{(3)} - r^{-2} k'' \left(2\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 9 \right) - r^{-3} k' \left(2\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 15 \right) \right. \\
 &\quad \left. + r^{-4} k \left(\left(\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} \right)^2 + 16\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 28 \right) \right) \tau \boxtimes dr \\
 &\quad - 12r^{-4} k \tilde{\underline{d}} \tilde{\underline{\delta}} \tau \boxtimes dr \\
 &\quad + 8 \left(-r^{-2} k'' - r^{-3} k' + r^{-4} k \left(\tilde{\underline{\nabla}}^* \tilde{\underline{\nabla}} + 6 \right) \right) \tilde{\underline{\delta}} \tau dr \otimes dr
 \end{aligned}$$

Proof. We must start from the expression in (2.14) and apply (2.13), (2.14), and (2.16) to the individual terms. First of all it is clear that if τ satisfies (2.5) and (2.6) then so does $\tilde{\underline{\nabla}}^{\text{sym}} \tau$, thus we may apply

(2.13) to conclude

$$\begin{aligned}
 (2.27) \quad & (\nabla^* \nabla) \left(-2r^{-2} k \underline{\widetilde{\nabla}}^{\text{sym}} \tau \right) \\
 &= \left(2r^{-2} k'' - 2r^{-3} k' - 2r^{-4} k \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 2 \right) \right) \left(\underline{\widetilde{\nabla}}^{\text{sym}} \tau \right) \\
 &\quad + 4r^{-4} k \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} - \underline{\widetilde{d}} \underline{\widetilde{\delta}} - 2 \right) \tau \boxtimes dr + 8r^{-4} k \underline{\widetilde{\delta}} \tau dr \otimes dr
 \end{aligned}$$

where we have used the equations

$$\underline{\widetilde{\delta}} \left(\underline{\widetilde{\nabla}}^{\text{sym}} \tau \right) = -\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} \tau + \underline{\widetilde{d}} \underline{\widetilde{\delta}} \tau + 2\tau, \quad \underline{\widetilde{\text{tr}}} \left(\underline{\widetilde{\nabla}}^{\text{sym}} \tau \right) = 2\underline{\widetilde{\delta}} \tau$$

from lemma 6.2. Now we apply (2.14) to conclude

$$\begin{aligned}
 (2.28) \quad & (\nabla^* \nabla) (-k'' \tau \boxtimes dr) \\
 &= 2r^{-2} k'' \underline{\widetilde{\nabla}}^{\text{sym}} \tau \\
 &\quad + \left\{ \left[k^{(4)} + 3r^{-1} k^{(3)} - r^{-2} k'' \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 6 \right) \right] \tau \right\} \boxtimes dr \\
 &\quad - 4r^{-2} k'' \underline{\widetilde{\delta}} \tau dr \otimes dr
 \end{aligned}$$

And similarly

$$\begin{aligned}
 (2.29) \quad & (\nabla^* \nabla) (-3r^{-1} k' \tau \boxtimes dr) \\
 &= 6r^{-3} k' \underline{\widetilde{\nabla}}^{\text{sym}} \tau \\
 &\quad + 3 \left\{ \left[r^{-1} k^{(3)} + r^{-2} k'' - r^{-3} k' \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 7 \right) \right] \tau \right\} \boxtimes dr \\
 &\quad - 12r^{-3} k' \underline{\widetilde{\delta}} \tau dr \otimes dr
 \end{aligned}$$

And again

$$\begin{aligned}
 (2.30) \quad & (\nabla^* \nabla) \left(r^{-2} k \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 6 \right) \tau \right) \boxtimes dr \\
 &= -2r^{-4} k \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 8 \right) \underline{\widetilde{\nabla}}^{\text{sym}} \tau \\
 &\quad + \left\{ \left[-r^{-2} k'' + r^{-3} k' + r^{-4} k \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 6 \right) \right] \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 6 \right) \tau \right\} \boxtimes dr \\
 &\quad + 4r^{-4} k \left(\underline{\widetilde{\nabla}}^* \underline{\widetilde{\nabla}} + 4 \right) \underline{\widetilde{\delta}} \tau dr \otimes dr
 \end{aligned}$$

where we commuted derivatives using lemma 6.2. Finally using (2.16) with $l = r^{-2}k$ and $\phi = \underline{\tilde{\delta}}\tau$ we compute

$$\begin{aligned}
 (2.31) \quad & (\nabla^* \nabla) \left(4r^{-2}k\underline{\tilde{\delta}}\tau \right) dr \otimes dr \\
 &= -8r^{-4}k\underline{\tilde{\delta}}\tau\underline{\tilde{g}} - 8r^{-4}k\underline{\tilde{d}}\underline{\tilde{\delta}}\tau \boxtimes dr \\
 &\quad + 4 \left[-r^{-2}k'' + r^{-3}k' + r^{-4}k \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 6 \right) \right] \left(\underline{\tilde{\delta}}\tau \right) dr \otimes dr
 \end{aligned}$$

Collecting together (2.27) - (2.31) gives the result. \square

Proposition: 2.7. *Consider $l(r)\phi(x)dr \otimes dr$. Then*

$$\begin{aligned}
 (2.32) \quad & (\nabla^* \nabla)^2 (l\phi dr \otimes dr) \\
 &= 4 \left(r^{-2}l'' + r^{-3}l' - r^{-4}l \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 4 \right) \right) (\phi\underline{\tilde{g}}) \\
 &\quad + 4r^{-4}l\underline{\tilde{\nabla}}^{\text{sym}}\underline{\tilde{d}}\phi \\
 &\quad + 4 \left(r^{-2}l'' + r^{-3}l' - r^{-4}l \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 8 \right) \right) \underline{\tilde{d}}\phi \boxtimes dr \\
 &\quad + \left(l^{(4)} + 6r^{-1}l^{(3)} - r^{-2}l'' \left(2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 9 \right) - r^{-3}l' \left(2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 15 \right) \right. \\
 &\quad \left. + r^{-4}l \left(\left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \right)^2 + 20\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 48 \right) \right) \phi dr \otimes dr
 \end{aligned}$$

Proof. We need to compute the action of the Laplacian of the right hand side of (2.16). First of all it is clear that $\phi\underline{\tilde{g}}$ satisfies (2.5) and (2.6) thus we may apply (2.13) to conclude

$$\begin{aligned}
 (2.33) \quad & (\nabla^* \nabla) \left(-2r^{-2}l\phi\underline{\tilde{g}} \right) \\
 &= \left(2r^{-2}l'' - 2r^{-3}l' - 2r^{-4}l \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 2 \right) \right) (\phi\underline{\tilde{g}}) \\
 &\quad - 4r^{-4}l\underline{\tilde{d}}\phi \boxtimes dr + 12r^{-4}l\phi dr \otimes dr
 \end{aligned}$$

where we used that $\underline{\tilde{\delta}}\underline{\tilde{g}} = 0$ and $\text{tr}\underline{\tilde{g}} = 3$. Next we note that $\underline{\tilde{d}}\phi$ satisfies (2.5) and (2.6) so that we can apply (2.14) to conclude

$$\begin{aligned}
 (2.34) \quad & (\nabla^* \nabla) \left(-2r^{-2}l \right) \underline{\tilde{d}}\phi \boxtimes dr \\
 &= 4r^{-4}l\underline{\tilde{\nabla}}^{\text{sym}}\underline{\tilde{d}}\phi \\
 &\quad + \left\{ \left[2r^{-2}l'' - 2r^{-3}l' - 2r^{-4}l \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 6 \right) \right] \underline{\tilde{d}}\phi \right\} \boxtimes dr \\
 &\quad + 8r^{-4}l\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \phi dr \otimes dr
 \end{aligned}$$

Where we have used the equation $\widetilde{\nabla}^* \widetilde{\nabla} = -\widetilde{\delta} \widetilde{d}$ Next using (2.16) we compute

$$\begin{aligned}
 (2.35) \quad & (\nabla^* \nabla) (-l'' \phi dr \otimes dr) \\
 &= 2r^{-2} l'' \phi \widetilde{g} + 2r^{-2} l'' \widetilde{d} \phi \boxtimes dr \\
 &\quad + \left(l^{(4)} + 3r^{-1} l^{(3)} - r^{-2} l'' \left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \right) \phi dr \otimes dr
 \end{aligned}$$

And similarly

$$\begin{aligned}
 (2.36) \quad & (\nabla^* \nabla) (-3r^{-1} l' \phi dr \otimes dr) \\
 &= 6r^{-3} l' \phi \widetilde{g} + 6r^{-3} l' \widetilde{d} \phi \boxtimes dr \\
 &\quad + 3 \left(r^{-1} l^{(3)} + r^{-2} l'' - r^{-3} l' \left(\widetilde{\nabla}^* \widetilde{\nabla} + 7 \right) \right) \phi dr \otimes dr
 \end{aligned}$$

Finally we apply (2.16) to conclude

$$\begin{aligned}
 (2.37) \quad & (\nabla^* \nabla) \left(r^{-2} l \left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \phi dr \otimes dr \right) \\
 &= -2r^{-4} l \left[\left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \phi \right] \widetilde{g} - 2r^{-4} l \left[\left(\widetilde{\nabla}^* \widetilde{\nabla} + 8 \right) \widetilde{d} \phi \right] \boxtimes dr \\
 &\quad + \left(-r^{-2} l'' + r^{-3} l' + r^{-4} l \left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \right) \left[\left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \phi \right] dr \otimes dr
 \end{aligned}$$

Collecting together (2.33) - (2.37) gives the result. \square

3. THE LINEARIZED EQUATION

In this section we derive the equation for the linearization of the Bach-flat condition at a flat metric. The equation is of course not strictly elliptic due to the diffeomorphism invariance of the Bach-flat equation. After restricting to the case where the variation of the metric is trace-free and divergence-free, we are reduced to the Biharmonic equation. We can then use the decomposition of the Bilaplacian computed in the previous section to classify solutions to this equation and moreover compute the exact decay rates of solutions, which is the key ingredient of our removal of singularities result.

Proposition: 3.1. *Given g a flat metric, the linearization of the Bach-flat equation at g is equivalent to the equation*

$$\begin{aligned}
 (3.1) \quad & 0 = \Delta^2 h - 2\Delta \delta^* \delta h \\
 & \quad + \frac{1}{3} \left(\nabla^2 \Delta \operatorname{tr}_g h - (\Delta^2 \operatorname{tr}_g h) g + (\Delta \delta^2 h) g + 2\nabla^2 \delta^2 h \right)
 \end{aligned}$$

Proof. Let $g(s)$ be a family of metrics such that $g(0) = g$ and $\frac{\partial}{\partial s}g(s)|_{s=0} = h$. First of all since the metric is flat it is clear that

$$\left(\frac{1}{2}R^{kl}W_{ikjl}\right)' \cdot h = 0$$

As a consequence of the Bianchi identity, in dimension 4 we have the equation

$$\nabla^k \nabla^l W_{ikjl} = \Delta A_{ij} - \nabla^k \nabla_i A_{jk}$$

where A_{ij} are the components of the Weyl-Schouten tensor

$$A = \frac{1}{n-2} \left(\text{Rc} - \frac{1}{2(n-1)} R \cdot g \right)$$

A further application of the Bianchi identity then yields

$$\nabla^k \nabla^l W_{ikjl} = \frac{1}{2} \left(\Delta \text{Rc}_{ij} - \frac{1}{6} \Delta R g_{ij} - \frac{1}{3} \nabla_i \nabla_j R \right)$$

First we compute the linearization

$$(\Delta \text{Rc})' \cdot h = -\frac{1}{2} \Delta (\Delta h + \nabla^2 \text{tr } h - 2\delta^* \delta h)$$

and next

$$\left(-\frac{1}{6} \Delta R g\right)' \cdot h = \frac{1}{6} \Delta (\Delta \text{tr } h - \delta \delta h) g$$

and finally

$$\left(-\frac{1}{3} \nabla^2 R\right)' \cdot h = \frac{1}{3} \nabla^2 (\Delta \text{tr } h - \delta \delta h)$$

The result follows. \square

Now, by a diffeomorphism gauge-fixing procedure which we will use in the applications, we may assume that h comes to us almost divergence-free. In fact we cannot prescribe that h is exactly divergence free. We will address this difficulty later in this section. Also, because of the conformal invariance of the Bach-flat equation, the trace of the right hand side of (3.1) vanishes. Thus there is no a-priori separate equation that the trace of h must satisfy. However, using this conformal invariance in the applications, we can assume that our Bach-flat metric has constant scalar curvature, which in the asymptotically flat regime corresponds to the equation

$$(3.2) \quad \Delta \text{tr } h = 0$$

As it turns out we can now remove the trace entirely by adding a term $L_X g$ for some vector field X . This corresponds to solving for a vector field X satisfying

$$(3.3) \quad \delta L_X g = 0$$

to preserve the divergence-free condition and

$$(3.4) \quad \text{tr}(L_X g) = 2\delta X = u$$

where u is a prescribed harmonic function. It is shown in [1] lines 4.11-4.20 that on a Ricci-flat cone one can indeed solve for this X . Thus a divergence-free solution to (3.1) satisfying (3.2) can be written as $L_X g + h$ where

$$(3.5) \quad \begin{aligned} 0 &= \Delta^2 h \\ 0 &= \text{tr}_g h \\ 0 &= \delta h \end{aligned}$$

We now proceed to analyse solutions to (3.5). Let $h = f(r)B + k(r)\tau \boxtimes dr + l(r)\phi(x)dr \otimes dr$. First of all the equation $\text{tr} h = 0$ reduces to the equation

$$(3.6) \quad f\widetilde{\text{tr}}(B) + l\phi = 0$$

Next, using lemma 2.4 and (3.6) the equation $\delta h = 0$ reduces to the equations

$$(3.7) \quad r^{-1}f\widetilde{\delta}(B) + (k' + 4r^{-1}k)\tau = 0$$

$$(3.8) \quad (l' + 4r^{-1}l)\phi + r^{-1}k(\widetilde{\delta}\tau) = 0$$

Now we would like to write down and simplify the vertical component of equation (3.5). Using propositions 2.20, 2.26 and 2.32 we see that this equation gives

$$(3.9) \quad \begin{aligned} 0 &= \left(l^{(4)} + 6r^{-1}l^{(3)} - r^{-2}l'' \left(2\widetilde{\nabla}^* \widetilde{\nabla} + 9 \right) - r^{-3}l' \left(2\widetilde{\nabla}^* \widetilde{\nabla} + 15 \right) \right. \\ &\quad \left. + r^{-4}l \left(\left(\widetilde{\nabla}^* \widetilde{\nabla} \right)^2 + 20\widetilde{\nabla}^* \widetilde{\nabla} + 48 \right) \right) \phi \\ &\quad + 8 \left(-r^{-2}k'' - r^{-3}k' + r^{-4}k \left(\widetilde{\nabla}^* \widetilde{\nabla} + 6 \right) \right) \widetilde{\delta}\tau \\ &\quad + 4 \left(r^{-2}f'' + r^{-3}f' - r^{-4}f \left(\widetilde{\nabla}^* \widetilde{\nabla} + 4 \right) \right) \widetilde{\text{tr}}(B) \\ &\quad + 8r^{-4}f\widetilde{\delta}\widetilde{\delta}(B) \end{aligned}$$

Now we want to use equations (3.6), (3.7) and (3.8) to simplify this further. First of all using (3.6) it is clear that

$$(3.10) \quad \begin{aligned} & 4 \left(r^{-2} f'' + r^{-3} f' - r^{-4} f \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 4 \right) \right) \widetilde{\text{tr}}(B) \\ &= -4 \left(r^{-2} l'' + r^{-3} l' - r^{-4} l \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 4 \right) \right) \phi \end{aligned}$$

Next using (3.8) we compute

$$(3.11) \quad \begin{aligned} & 8 \left(-r^{-2} k'' - r^{-3} k' + r^{-4} k \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 6 \right) \right) \widetilde{\underline{\delta}} \tau \\ &= 8 \left(r^{-1} l^{(3)} + 7r^{-2} l'' - r^{-3} l' \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 1 \right) - 4r^{-4} l \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 6 \right) \right) \phi \end{aligned}$$

Finally using both (3.7) and (3.8) we compute

$$(3.12) \quad \begin{aligned} 8r^{-4} f \widetilde{\underline{\delta}} \widetilde{\underline{\delta}}(B) &= -8 \left(r^{-3} k' + 4r^{-4} k \right) \widetilde{\underline{\delta}} \tau \\ &= 8 \left(r^{-2} l'' + 9r^{-3} l' + 16r^{-4} l \right) \phi \end{aligned}$$

Thus plugging (3.10) - (3.12) into (3.9) gives the equation

$$(3.13) \quad \begin{aligned} 0 &= \left(l^{(4)} + 14r^{-1} l^{(3)} + r^{-2} l'' \left(-2\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 51 \right) \right. \\ &\quad \left. + r^{-3} l' \left(-10\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 45 \right) + r^{-4} l \left(\left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} \right)^2 - 8\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} \right) \right) \phi \end{aligned}$$

We would now like to do the same for the cross component of (3.5). First of all from propositions 2.20, 2.26 and 2.32 we see that this is the equation

$$(3.14) \quad \begin{aligned} 0 &= 4 \left[-r^{-2} f'' - r^{-3} f' + r^{-4} f \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 2 \right) \right] \widetilde{\underline{\delta}}(B) + 8r^{-4} f \widetilde{\underline{d}} \widetilde{\text{tr}}(B) \\ &\quad + \left(k^{(4)} + 6r^{-1} k^{(3)} - r^{-2} k'' \left(2\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 9 \right) - r^{-3} k' \left(2\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 15 \right) \right. \\ &\quad \left. + r^{-4} k \left(\left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} \right)^2 + 16\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 28 \right) \right) \tau - 12r^{-4} k \widetilde{\underline{d}} \widetilde{\underline{\delta}} \tau \\ &\quad + 4 \left(r^{-2} l'' + r^{-3} l' - r^{-4} l \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 8 \right) \right) \widetilde{\underline{d}} \phi \end{aligned}$$

Now, using (3.7) we see that

$$(3.15) \quad \begin{aligned} & 4 \left[-r^{-2} f'' - r^{-3} f' + r^{-4} f \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 2 \right) \right] \widetilde{\underline{\delta}}(B) \\ &= 4 \left(r^{-1} k^{(3)} + 7r^{-2} k'' + r^{-3} k' \left(-\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 3 \right) \right. \\ &\quad \left. - 4r^{-4} k \left(\widetilde{\underline{\nabla}}^* \widetilde{\underline{\nabla}} + 2 \right) \right) \tau \end{aligned}$$

Next using (3.6) we make the simplification

$$(3.16) \quad 8r^{-4}f\widetilde{d}\widetilde{\text{tr}}(B) = -8r^{-4}l\widetilde{d}\phi$$

And using (3.8) we simplify

$$(3.17) \quad -12r^{-4}k\widetilde{d}\widetilde{\delta}\tau = 12(r^{-3}l' + 4r^{-4}l)\widetilde{d}\phi$$

Plugging (3.15) - (3.17) into (3.14) gives

$$(3.18) \quad \begin{aligned} 0 = & \left(k^{(4)} + 10r^{-1}k^{(3)} + r^{-2}k'' \left(-2\widetilde{\nabla}^* \widetilde{\nabla} + 19 \right) \right. \\ & - k' \left(6\widetilde{\nabla}^* \widetilde{\nabla} + 3 \right) + r^{-4}k \left(\left(\widetilde{\nabla}^* \widetilde{\nabla} \right)^2 - 4 \right) \Big) \tau \\ & + 4 \left(r^{-2}l'' + 4r^{-3}l' + r^{-4}l \left(-\widetilde{\nabla}^* \widetilde{\nabla} + 2 \right) \right) \widetilde{d}\phi \end{aligned}$$

Finally we need to simplify the horizontal component of (3.5). Again using propositions 2.20, 2.26 and 2.32 we get the equation

$$(3.19) \quad \begin{aligned} 0 = & \left(f^{(4)} + 6r^{-1}f^{(3)} - r^{-2}f'' \left(1 + 2\widetilde{\nabla}^* \widetilde{\nabla} \right) \right. \\ & - r^{-3}f' \left(2\widetilde{\nabla}^* \widetilde{\nabla} + 7 \right) + r^{-4}f \left(\widetilde{\nabla}^* \widetilde{\nabla} + 2 \right)^2 \Big) B \\ & - 4r^{-4}f\widetilde{\nabla}^{\text{sym}}\widetilde{\delta}(B) + 4r^{-4}f\widetilde{\text{tr}}(B)\widetilde{g} \\ & + 4 \left(r^{-2}k'' + r^{-3}k' - r^{-4}k \left(\widetilde{\nabla}^* \widetilde{\nabla} + 5 \right) \right) \left(\widetilde{\nabla}^{\text{sym}}\tau \right) - 8r^{-4}k\widetilde{\delta}\tau\widetilde{g} \\ & + 4 \left(r^{-2}l'' + r^{-3}l' - r^{-4}l \left(\widetilde{\nabla}^* \widetilde{\nabla} + 4 \right) \right) (\phi\widetilde{g}) \\ & + 4r^{-4}l\widetilde{\nabla}^{\text{sym}}\widetilde{d}\phi \end{aligned}$$

First using (3.7) we simplify

$$(3.20) \quad -4r^{-4}f\widetilde{\nabla}^{\text{sym}}\widetilde{\delta}(B) = 4(r^{-3}k' + 4r^{-4}k)\widetilde{\nabla}^{\text{sym}}\tau$$

Using (3.6) we simplify

$$(3.21) \quad 4r^{-4}f\widetilde{\text{tr}}(B)\widetilde{g} = -4r^{-4}l\phi\widetilde{g}$$

Next from (3.8) we simplify

$$(3.22) \quad -8r^{-4}k\widetilde{\delta}\tau\widetilde{g} = 8(r^{-3}l' + 4r^{-4}l)\phi\widetilde{g}$$

Plugging (3.20) - (3.22) into (3.19) gives

$$\begin{aligned}
 (3.23) \quad 0 = & \left(f^{(4)} + 6r^{-1}f^{(3)} - r^{-2}f'' \left(1 + 2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \right) \right. \\
 & \left. - r^{-3}f' \left(2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 7 \right) + r^{-4}f \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 2 \right)^2 \right) B \\
 & + 4 \left(r^{-2}k'' + 2r^{-3}k' - r^{-4}k \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 1 \right) \right) \left(\underline{\tilde{\nabla}}^{\text{sym}} \tau \right) \\
 & + 4 \left(r^{-2}l'' + 3r^{-3}l' + r^{-4}l \left(-\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 3 \right) \right) \phi \underline{\tilde{g}} + 4r^{-4}l \underline{\tilde{\nabla}}^{\text{sym}} \underline{\tilde{d}}\phi
 \end{aligned}$$

Lemma: 3.2. *The system of equations (3.5) is equivalent to*

$$\begin{aligned}
 (3.24) \quad 0 = & \left(f^{(4)} + 6r^{-1}f^{(3)} - r^{-2}f'' \left(1 + 2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \right) \right. \\
 & \left. - r^{-3}f' \left(2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 7 \right) + r^{-4}f \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 2 \right)^2 \right) B \\
 & + 4 \left(r^{-2}k'' + 2r^{-3}k' - r^{-4}k \left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 1 \right) \right) \left(\underline{\tilde{\nabla}}^{\text{sym}} \tau \right) \\
 & + 4 \left(r^{-2}l'' + 3r^{-3}l' + r^{-4}l \left(-\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 3 \right) \right) \phi \underline{\tilde{g}} + 4r^{-4}l \underline{\tilde{\nabla}}^{\text{sym}} \underline{\tilde{d}}\phi
 \end{aligned}$$

$$\begin{aligned}
 (3.25) \quad 0 = & \left(k^{(4)} + 10r^{-1}k^{(3)} + r^{-2}k'' \left(-2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 19 \right) \right. \\
 & \left. - r^{-3}k' \left(6\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 3 \right) + r^{-4}k \left(\left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \right)^2 - 4 \right) \right) \tau \\
 & + 4 \left(r^{-2}l'' + 4r^{-3}l' + r^{-4}l \left(-\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 2 \right) \right) \underline{\tilde{d}}\phi
 \end{aligned}$$

$$\begin{aligned}
 (3.26) \quad 0 = & \left(l^{(4)} + 14r^{-1}l^{(3)} + r^{-2}l'' \left(-2\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 51 \right) \right. \\
 & \left. + r^{-3}l' \left(-10\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} + 45 \right) + r^{-4}l \left(\left(\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \right)^2 - 8\underline{\tilde{\nabla}}^* \underline{\tilde{\nabla}} \right) \right) \phi
 \end{aligned}$$

$$(3.27) \quad 0 = r^{-1}f \underline{\tilde{d}}(B) + (k' + 4r^{-1}k) \tau$$

$$(3.28) \quad 0 = (l' + 4r^{-1}l) \phi + r^{-1}k \left(\underline{\tilde{d}}\tau \right)$$

We now analyse all solutions to this system of equations. We notice that for any solution, l satisfies the determined ODE in (3.26). Thus we can compute the solutions to (3.26) in terms of the eigenvalues of the Laplacian on S^3 acting on functions. Once these are determined we notice that the equations (3.27) and (3.28) show that both f and k will

have the same decay rate as l . Once this is done we restrict to the case where $l = 0$. Here we notice that k satisfies the now determined ODE in (3.25). We can classify these solutions in terms of the eigenvalues of the Laplacian on S^3 acting on one-forms. Once again we can use equation (3.27) to conclude that f has the same decay rate as k . Finally we restrict to the case where $l = k = 0$. Then it is clear that f satisfies the determined ODE in (3.24), and classify solutions in terms of the eigenvalues of the Laplacian on S^3 acting on traceless symmetric two-tensors. We now make this rigorous in a series of lemmas.

Lemma: 3.3. *The solutions to (3.5) satisfying $l = k = 0$ are*

$$h = \sum_i r^{b_i} B_i$$

where either $b_i > 0$ or

$$b_i \leq -2$$

Proof. Since $k = l = 0$ we have $\tilde{\text{tr}} B = \tilde{\delta} B = 0$. Say λ is an eigenvalue for the Laplacian of S^3 acting on traceless symmetric two-tensors, and suppose B_λ is in the eigenspace of λ . Equation (3.24) then reduces to an ODE with solutions

$$(3.29) \quad r^{\pm 1 \pm \sqrt{\lambda+3}} B_\lambda.$$

From lemma 6.1 we know that the smallest eigenvalue is 6, so the result follows. \square

Lemma: 3.4. *The solutions to (3.5) satisfying $l = 0$ are*

$$h = L_X g_0 + \sum_i r^{b_i} (B_i + \tau_i \boxtimes dr)$$

where either $b_i > 0$ or

$$b_i \leq -2$$

Proof. Note that since $l = 0$ equation (3.28) implies in that $\tilde{\delta} \tau = 0$. The eigenvalues for the Laplacian acting on 1-forms on S^3 are given by [5]

$$(3.30) \quad a_j := (j+1)(j+3)$$

Say τ_j is in the eigenspace of a_j . Then (3.25) reduces to an ODE with solutions

$$(3.31) \quad r^{-2 \pm \sqrt{a_j+2}} \tau_j$$

We point out that the decay rates are all less than -2 . Also, since the decay rate -4 does not occur, the expression $k' + 4r^{-1}k$ never vanishes,

thus by equation (3.27) we see that $f = cr^{-2\pm\sqrt{a_j+2}}$, and the result follows. \square

Lemma: 3.5. *The solutions to (3.5) satisfying $l \neq 0$ are*

$$h = L_X g_0 + \sum_i r^{b_i} (B_i + \tau_i \boxtimes dr + \phi dr \otimes dr)$$

where either $b_i > 0$ or

$$b_i \leq -2$$

Proof. The eigenvalues of the Laplacian acting on functions on S^3 are

$$(3.32) \quad a_j := j(j+2)$$

Suppose ϕ_j is in the eigenspace of a_j . Equation (3.26) reduces to an ODE with solutions

$$(3.33) \quad l(r) = r^{-2\pm 1\pm\sqrt{a_j+1}}$$

In the case where $l(r) = r^{-3\pm\sqrt{a_j+1}}$, the solution to the whole system is given by a Lie derivative term, specifically $L_X g_0$ where

$$X_* = r^{-1\pm\sqrt{a_j+1}} d\phi_j + (-1 \pm \sqrt{a_j+1}) r^{-2\pm\sqrt{a_j+1}} \phi_j dr.$$

For the solutions $l(r) = r^{-1\pm\sqrt{a_j+1}}$ we note that for $j > 0$ the decay rates are always less than or equal to -2 . In the case $j = 0$, i.e. where $\phi = c$, consider the radially parallel solution, i.e. where $l(r) = c$. Using equation (3.28) we see that k is also a nonzero constant. Using this equation (3.26) reduces to $\left(\left(\widetilde{\nabla}^* \widetilde{\nabla} \right)^2 - 4 \right) \tau = 0$ so we see that this solution does not in fact occur since 2 is not an eigenvalue of the Laplacian on S^3 acting on one-forms.

In any of the cases above, using equation (3.6) we see that f will have the same decay rate as l . Given this, equation (3.27) implies that k must also have the same decay rate. \square

Together the above lemmas and earlier calculations prove the following proposition.

Proposition: 3.6. *On a flat cone, solutions of (3.1) satisfying $\delta h = 0$ and $\Delta \operatorname{tr} h = 0$ can be written uniquely as a sum*

$$(3.34) \quad h = L_X g_0 + \sum_i r^{b_i} (B_i + \tau_i \boxtimes dr + \phi dr \otimes dr)$$

where either $b_i > 0$ or

$$(3.35) \quad b_i \leq -2$$

We hasten to point out that we will *not* in fact be able to guarantee the divergence-free condition. This is due to the presence of certain eigenvalues of the Laplace-Beltrami operator acting on one-forms. Thus we follow the technique used in [1] and consider a modified divergence-free condition. In particular, fix $t \neq 0$ with $|t|$ very small. Let

$$(3.36) \quad \delta_t = \delta - ti_{r^{-1}} \frac{\partial}{\partial r}$$

We will be able to prescribe that $\delta_t h = 0$ for t arbitrarily small. Given this, we define the following modified equation

$$(3.37) \quad \begin{aligned} 0 &= \Delta^2 h - 2t\Delta\delta^* i_{r^{-1}} \frac{\partial}{\partial r} h \\ &\quad + \frac{1}{3} \left(\nabla^2 \Delta \operatorname{tr}_g h - (\Delta^2 \operatorname{tr}_g h) g + t \left(\Delta \delta i_{r^{-1}} \frac{\partial}{\partial r} h \right) g + 2t \nabla^2 \delta i_{r^{-1}} \frac{\partial}{\partial r} h \right) \\ &:= \mathcal{P}(h) \end{aligned}$$

Following the analysis of proposition 3.6 we can write the solutions satisfying $\Delta \operatorname{tr} h = 0$ as growth and decay solutions where the rates are perturbations of those calculated above. In particular we can write

$$(3.38) \quad h = r^{\beta_i} T_i$$

where T_i is some symmetric bilinear form and $\{T_i\}$ are orthonormal with respect to the inner product

$$(3.39) \quad \langle\langle h_1, h_2 \rangle\rangle = r^{-3} \int_{(r, S^3)} \langle h_1, h_2 \rangle dV_{S^3}$$

We also have the following corollary.

Corollary: 3.7. *For t sufficiently small, there are no radially parallel solutions of (3.37) satisfying $\Delta \operatorname{tr} h = 0$.*

Proof. We note that the radially parallel solutions found above all had a dr component, and so in the perturbed equation are no longer radially parallel. Thus the only possibility would be $f(r)B$ where $f(r)$ is a constant function and B is trace and divergence-free. However, no such solution occurs according to our analysis above. Thus the corollary follows. \square

So, according to the above results we can decompose any solution satisfying $\Delta \operatorname{tr} h = 0$ as

$$(3.40) \quad h = h_{\uparrow} + h_{\downarrow}$$

where h_\uparrow are the solutions with positive growth rate and likewise h_\downarrow are solutions which decay in r . Now let

$$(3.41) \quad \beta = \min_i |\beta_i| > 0$$

We can now state the first of our decay estimates. We start with some notation. For a fixed annulus $A_{a,b}(p)$ we have the norm

$$(3.42) \quad |||h|||_{a,b} = \int_a^b ||h||^2 r^{-1} dr$$

where the norm $|| \cdot ||$ is defined in 3.39. This norm is defined so that if $w = a^{-2} \psi_a^*(h)$, where ψ_a is the natural scaling map $\psi_a(x, r) = (x, ar)$, then

$$(3.43) \quad |||h|||_{a,La} = |||w|||_{1,L}$$

Corollary: 3.8. *Given $0 < \beta' < \beta$ there exists l such that for all $a > 0$ and $L \geq l$*

$$(3.44) \quad |||h_\uparrow|||_{La,L^2a} \geq L^{\beta'} |||h_\uparrow|||_{a,La}$$

$$(3.45) \quad |||h_\downarrow|||_{La,L^2a} \leq L^{-\beta'} |||h_\downarrow|||_{a,La}$$

Lemma: 3.9. *Given $0 < \beta' < \beta$ there exists L such that if k is a solution to (3.37) satisfying $\delta_t k = 0$ and $\Delta \operatorname{tr} k = 0$ then if*

$$(3.46) \quad |||k|||_{L,L^2} \geq L^{\beta'} |||k|||_{1,L}$$

then

$$(3.47) \quad |||k|||_{L^2,L^3} \geq L^{\beta'} |||k|||_{L,L^2}$$

and if

$$(3.48) \quad |||k|||_{L^2,L^3} \leq L^{-\beta'} |||k|||_{L,L^2}$$

then

$$(3.49) \quad |||k|||_{L,L^2} \leq L^{-\beta'} |||k|||_{1,L}$$

4. ASYMPTOTIC CURVATURE DECAY

In this section we give the proof of theorem 1.1. The proof will follow the techniques used in [1]. We first recall the definition of certain norms introduced in [1]. Let A_u denote the natural action of the scaling ψ_u on tensors of type (p, q) , i.e.

$$(4.1) \quad A_u = (\psi_{u^{-1}})_* \otimes \cdots \otimes (\psi_{u^{-1}})_* \otimes \psi_u^* \otimes \cdots \otimes \psi_u^*$$

Given T a tensor of type (p, q) we have

$$|T(u, x)|_{k,\alpha;0} = |u^{p-q} A_u T(1, x)|_{k,\alpha;0}$$

where the norm on the right is the $C^{k,\alpha}$ -norm with respect to g_0 at the point $(1, x)$. Using this we define

$$(4.2) \quad |T|_{k,\alpha;0} = \sup_{(u,x)} |T(u, x)|_{k,\alpha;0}$$

And more generally

$$(4.3) \quad |T|_{k,\alpha;l} = |r^{-l}T|_{k,\alpha;0}$$

Given T a tensor of type (p, q) we write $T \in \mathcal{T}_{k,\alpha;l}^{p,q}$ if $|T|_{k,\alpha;l} < \infty$.

We will make use of a gauge-fixing theorem ([1] Theorem 3.1). For a given cone with cone-point p , let $A_{c,d} = \{(r, x) | c < r < d\}$.

Proposition: 4.1. *Fix $t \neq 0$. There exists $\kappa(t, k)$ such that if $(C(N^{n-1}), g_0)$ is a Ricci-flat cone and g is a metric on $A_{c,d}(p) \subset C(N^{n-1})$ where $d/c \geq 2$ such that*

$$(4.4) \quad |g - g_0|_{k,\alpha;0} < \kappa(t, k)$$

then there exists a diffeomorphism $\phi : A_{c,d}(p) \rightarrow A_{c,d}(p)$ such that

$$(4.5) \quad \phi^*g \in \mathcal{T}_{k,\alpha;0}^{0,2}$$

and

$$(4.6) \quad \delta_t(\phi^*g - g_0) = 0$$

Since we have assumed our given Bach-flat metric is asymptotically flat, we can use this theorem to construct gauges relative to the flat metric on the cone in the asymptotic regime. Specifically, let g_0 be the flat metric on \mathbb{R}^4 and let g be a given Bach-flat metric on this cone. Consider a fixed annulus $A_{c,d}$ and let

$$(4.7) \quad h = \phi^*g - g_0$$

where ϕ is from proposition 4.1 with respect to this fixed annulus. Then we note that h satisfies

$$(4.8) \quad \delta_t h = 0$$

and further

$$(4.9) \quad B_{g_0+h} - B_{g_0} = 0$$

which is a nonlinear elliptic equation on h which we can think of as a perturbation of the linearized deformation equation for h small. The following lemma makes this precise.

Lemma: 4.2.

$$(4.10) \quad B_{g_0+h} - B_{g_0} = \mathcal{P}(h) + F(h, g_0)$$

where

(4.11)

$$|F|_{k,\alpha;0} \leq C \sum_{p+q+r+s=4} |\nabla^p h|_{k-p,\alpha;-p} |\nabla^q h|_{k-q,\alpha;-q} |\nabla^r h|_{k-r,\alpha;-r} |\nabla^s h|_{k-s,\alpha;-s}$$

Also, since the metric g is scalar-flat, we get the equation

$$(4.12) \quad R_{g_0+h} - R_{g_0} = 0$$

from which we derive a similar lemma.

Lemma: 4.3.

$$(4.13) \quad R_{g_0+h} - R_{g_0} = -\Delta \operatorname{tr} h - t\delta \left(i_{r^{-1}} \frac{\partial}{\partial r} h \right) + F'(h, g_0)$$

where

$$(4.14) \quad |F'|_{k,\alpha;0} \leq C \sum_{p+q=2} |\nabla^p h|_{k-p,\alpha;-p} |\nabla^q h|_{k-q,\alpha;-q}$$

Lemma: 4.4. *There exists a small constant $\chi(\Lambda, K)$ such that if $|h|_{k,\alpha;0} \leq \chi$ then*

$$(4.15) \quad |h(2a)|_{k,\alpha;0} \leq c(\Lambda, K, k) |||h|||_{a,4a}$$

Proof. Using the scale-invariance of the respective norms, it suffices to consider the case $a = 1$. The result follows in this case by standard elliptic theory. \square

Now let π denote orthogonal projection onto $\ker(\mathcal{P}(h))$ with respect to the inner product defining the norm $||| \cdot |||_{La, L^2a}$. Similar to (3.40) we can write

$$(4.16) \quad \pi h = (\pi h)_\uparrow + (\pi h)_\downarrow$$

Proposition: 4.5. *There exists $\chi = \chi(\Lambda, K) > 0$ and L large so that if $|h|_{k,\alpha;0} < \chi$ then if*

$$(4.17) \quad |||h|||_{La, L^2a} \geq L^{\beta'} |||h|||_{a, La}$$

then

$$(4.18) \quad |||h|||_{L^2a, L^3a} \geq L^{\beta'} |||h|||_{La, L^2a}$$

and if

$$(4.19) \quad |||h|||_{L^2a, L^3a} \leq L^{-\beta'} |||h|||_{La, L^2a}$$

then

$$(4.20) \quad |||h|||_{La, L^2a} \leq L^{-\beta'} |||h|||_{a, La}$$

Moreover, at least one of (4.17) or (4.19) must hold.

Proof. Suppose one had a sequence of gauges ϕ_i and solutions h_i where $|h_i|_{k,\alpha;0} \rightarrow 0$ but none of the assertions of the proposition hold. By rescaling our solutions and using a compactness argument (see [1] lemma 5.22) we can produce a solution to (3.37) satisfying $\Delta \operatorname{tr} h = 0$ which contradicts lemma 3.9. \square

Proof of Theorem 1.1 We will adopt the notation used in this section. In particular choose χ and L as in proposition 4.5. Let g_0 denote the standard flat metric on \mathbb{R}^4 . Since we already know that the metric g is ALE of order 0, there exists $\Psi : \mathbb{R}^4 \rightarrow X$ so that

$$(4.21) \quad |\Psi^*g - g_0|_{k,\alpha;0} = o(r)$$

Now consider the sequence of annuli $A_{L^i a, L^{i+1} a}(0)$. Using (4.21) we may apply proposition 4.1 and choose a δ_t -free gauge for Ψ^*g with respect to g_0 , ϕ_i , on each annulus. By pushing forward this is equivalent to finding a sequence of flat metrics g_i where $\Psi^*g - g_i$ is δ_t -free with respect to g_i .

$$h_i := \Psi^*g - g_i$$

We note that

$$\lim_{i \rightarrow \infty} |||h_i|||_{L^i a, L^{i+1} a} = 0.$$

Indeed, if this were not the case then inductively applying proposition 4.5, and in particular using that (4.17) implies (4.5) we can contradict (4.21). Since for any of the flat metrics g_i , there are no radially parallel solutions to the linearized deformation equation, we can conclude that (4.20) holds for all i .

Also, by passing to a subsequence, we can assume that for some flat metric g_∞ we have

$$(4.22) \quad \lim_{j \rightarrow \infty} |g_j - g_\infty|_{k,\alpha';0} = 0 \quad (\alpha' < \alpha)$$

Using this together with (4.20) and lemma 4.4 we conclude that

$$(4.23) \quad |\Psi^*g - g_\infty|_{k,\alpha;0} \leq cr^{-\beta'}$$

This proves that g is ALE of order β' .

To prove the full claim we will use an inductive procedure choosing better and better gauges. First of all we point out that the decay estimate (4.23) in fact suffices to find a completely divergence-free global gauge ϕ (Remark 3.23 [1]). So, let $h = \phi^*g - g_0$ be divergence-free. First we get a better estimate on the trace of h . In particular, by lemma 4.3 we have

$$(4.24) \quad \Delta \operatorname{tr} h = F'(h)$$

where the estimate (4.23) implies

$$(4.25) \quad F'(h) \in \mathcal{T}_{k-2,\alpha;-2\beta-2}^{0,2}$$

Applying the Greens function we can conclude the existence of a function f so that

$$(4.26) \quad \Delta f = F'(h)$$

where

$$(4.27) \quad f \in \mathcal{T}_{k,\alpha;-2\beta}^{0,2}$$

In particular this implies the equation

$$(4.28) \quad \Delta(\operatorname{tr} h - f) = 0$$

As a harmonic function on \mathbb{R}^4 , we know that the decay rate is at least -2 , thus

$$(4.29) \quad \operatorname{tr} h - f \in \mathcal{T}_{k,\alpha;-2}^{0,2}$$

This argument already can be inducted to show

$$(4.30) \quad \operatorname{tr} h \in \mathcal{T}_{k,\alpha;2-\delta}^{0,2}$$

for any $\delta > 0$. We now show how to get the rest of the estimate. By lemma 4.2 we have

$$(4.31) \quad \mathcal{P}(h) = F(h)$$

where \mathcal{P} now refers to the case $t = 0$. Using (4.23) we conclude

$$(4.32) \quad F(h) \in \mathcal{T}_{k-4,\alpha;-4\beta-4}^{0,2}$$

Again using a Greens function, this decay is sufficient to conclude the existence of h_1 such that

$$(4.33) \quad \mathcal{P}h_1 = F(h)$$

where

$$(4.34) \quad h_1 \in \mathcal{T}_{k,\alpha;-4\beta}^{0,2}$$

Thus $\mathcal{P}(h - h_1) = 0$. Given our estimate on the trace of h we can assume that $h - h_1$ is trace-free. We note that \mathcal{P} acting on the traceless piece of $h - h_1$ also vanishes. Using proposition 3.6 we can write

$$(4.35) \quad h - h_1 = L_X g_0 + h_2$$

where

$$(4.36) \quad h_2 \in \mathcal{T}_{k,\alpha;-\min\{4\beta,2\}}^{0,2}$$

and also

$$(4.37) \quad X \in \mathcal{T}_{k+1,\alpha;1-\beta}^{1,0}$$

Let K_X be the diffeomorphism generated by taking the flow of X to time 1. It is clear from the above estimates that

$$(4.38) \quad K_X^*(g_0) - L_X g_0 \in \mathcal{T}_{k,\alpha; -\min\{4\beta, 2\}}^{0,2}$$

Furthermore putting together the above estimates we get

$$(4.39) \quad K_{-X}^* \phi^* g - g_0 \in \mathcal{T}_{k,\alpha; -\min\{4\beta, 2\}}^{0,2}$$

If $4\beta \geq 2$ we are done. If not, start over with $h = K_{-X}^* \phi^* g - g_0$ and proceed as above with β replaced by $4\beta - \epsilon$ for very small $\epsilon > 0$. It is clear that by induction the result follows. \square

Remark: The result holds assuming harmonic curvature instead of the Bach-flat condition. In particular, using the Bianchi identity equation (1.12) implies an equation of the form $\Delta \text{Rc} = \text{Rm} * \text{Rm}$. The linearized deformation equation is the same as the one analyzed in section 3, and so one can apply the argument in theorem 1.1 to conclude that the metric is in fact ALE of order τ for any $\tau < 2$.

5. PROOF OF THEOREM 1.2

Let (X, d, x) and $B(x, 1) \setminus \{x\}$ be as in the statement of the theorem. We already know by theorem 1.1 of [7] that $B(x, 1)$ has a C^0 -orbifold structure at x . Thus in particular, lifting to the universal cover we may suppose $x = 0$ and $B(x, 1) \subset D(0, 1)$ where D denotes the distance ball in the Euclidean metric. Let g be this metric which is smooth away from the origin and C^0 on $B(0, 1)$. Fix a constant $s < 1$, and consider a sequence of annuli in the metric g

$$A_i := A(s^{i+1}, s^i)$$

Now let

$$(5.1) \quad \begin{aligned} \Phi : \mathbb{R}^4 \setminus \{0\} &\rightarrow \mathbb{R}^4 \setminus \{0\} \\ \Phi(r, \theta) &= \left(\frac{1}{r}, \theta \right) \end{aligned}$$

be the inversion through the unit sphere and let $A'_i = \Phi^* A_i$, and $g' = \Phi^* g$. Let ρ_i denote the distance function of the metric g' on the annulus A'_i . Since we have assumed a bound on the Sobolev constant, the L^2 -norm of curvature, and that the first Betti number is bounded, we may apply the analysis of section 4 in [7]. In particular we may conclude that $\rho_i \rightarrow r$ as $i \rightarrow \infty$. In particular the annuli A'_i are approaching the standard annuli in \mathbb{R}^4 . Thus the asymptotic analysis of section 4 above applies so we conclude that this metric is ALE of order τ for any

$\tau < 2$. Pulling this estimate back to B using the spherical inversion gives

$$(5.2) \quad g = \delta_{ij} + \mathcal{O}(r^{2-\delta}).$$

For δ sufficiently small this implies that the metric has a $C^{1,\alpha}$ extension through the origin. Using results from [3] this implies the existence of a harmonic coordinate around the origin. A simple computation using the Bianchi identity shows that the Bach-flat condition implies that the curvature satisfies an equation of the form $\Delta \text{Rc} = \text{Rm} * \text{Rc}$. Thus we view the Bach-flat equation in harmonic coordinates as the system

$$(5.3) \quad \Delta \text{Rc} = \text{Rm} * \text{Rc}$$

$$(5.4) \quad \Delta g = \text{Rc} + Q(g, \partial g)$$

Using the curvature estimate $|\text{Rm}| = \mathcal{O}(r^{-\delta})$ it follows from equation (5.3) that $\text{Rc} \in W^{2,p}$ for any p . Using equation (5.4) it is clear that $g \in W^{3,p}$. This allows us to bootstrap and conclude that $g \in C^\infty$. \square

Remark: Again we point out that the result holds assuming harmonic curvature instead of the Bach-flat condition. The crucial step in the above argument is the curvature decay rate, and we mentioned after the proof of theorem 1.1 that metrics with harmonic curvature satisfy this estimate.

6. APPENDIX: ANALYSIS ON THE THREE SPHERE

Lemma: 6.1. *Let (S^3, g) be the round three-sphere. The smallest eigenvalue of the rough Laplacian acting on traceless divergence-free symmetric two-tensors is 6.*

Proof. Using the identification of S^3 with $SU(2)$ we consider the standard global left-invariant Milnor frame X^1, X^2, X^3 with structure constants

$$C_{23}^1 = C_{31}^2 = C_{12}^3 = -2$$

If e_1, e_2, e_3 denotes the corresponding coframe, we have the equation

$$(6.1) \quad \nabla_{X^i} e_j = -e_k \sigma(ijk)$$

where $\sigma(ijk)$ denotes the sign of the permutation (ijk) and is zero if any of i, j and k are equal. Using this global frame we may write any traceless symmetric two-tensor as

$$B_{ij} e_i \otimes e_j$$

and compute

$$\begin{aligned}\nabla^* \nabla B = & -B_{kl,ii} e_k \boxtimes e_l - 2B_{kl,i} (\nabla_{X^i} e_k \boxtimes e_l + e_k \boxtimes \nabla_{X^i} e_l) \\ & - B_{kl} (\nabla_{X_i} \nabla_{X_i} e_k \boxtimes e_l + 2\nabla_{X^i} e_k \boxtimes \nabla_{X^i} e_l + e_k \boxtimes \nabla_{X^i} \nabla_{X^i} e_l)\end{aligned}$$

Now a basic calculation using (6.1) shows that

$$\nabla_{X^i} \nabla_{X^i} e_k = -2e_k$$

and similarly, using that B is traceless it is easy to compute

$$2B_{kl} \nabla_{X^i} e_k \boxtimes \nabla_{X^i} e_l = -2B_{kl} e_k \boxtimes e_l$$

Using these calculations and the fact that the smallest eigenvalue must occur when the coefficient functions are constant, the result follows. \square

Lemma: 6.2. *Let (S^3, g) be the round three-sphere. Then given a function $\phi \in C^\infty(S^3)$*

$$(6.2) \quad d\nabla^* \nabla \phi = \nabla^* \nabla d\phi + 2d\phi$$

and

$$(6.3) \quad \delta(\nabla^* \nabla) d\phi = -(\nabla^* \nabla)^2 \phi + 2\nabla^* \nabla \phi$$

and given a one-form $\alpha \in \wedge^1 T^* S^3$

$$(6.4) \quad \delta(\nabla^* \nabla) \alpha = \nabla^* \nabla \delta \alpha - 2\delta \alpha$$

and

$$(6.5) \quad \delta(\nabla^* \nabla)^2 \alpha = (\nabla^* \nabla)^2 \delta \alpha - 4\nabla^* \nabla \delta \alpha + 4\delta \alpha$$

and

$$(6.6) \quad \delta L_{\alpha^*} g = -\nabla^* \nabla \alpha + \nabla \delta \alpha + 2\alpha$$

and

$$(6.7) \quad L_{\nabla^* \nabla} g = \nabla^* \nabla L_\tau g + 4L_\tau g - 2\delta \tau g$$

and given B a traceless symmetric two-tensor

$$(6.8) \quad \delta \nabla^* \nabla B = \nabla^* \nabla \delta B - 4\delta B + 2d \operatorname{tr} B$$

and

$$(6.9) \quad \delta(\nabla^* \nabla)^2 = (\nabla^* \nabla)^2 \delta B - 8\nabla^* \nabla \delta B + 16\delta B$$

Proof. We use the convention that $R_{ijij} > 0$, and more specifically choosing normal coordinates at one point we have

$$\begin{aligned}R_{ijkl} &= (g_{ik}g_{jl} - g_{jk}g_{il}) \\ R_{ik} &= 2g_{ik} \\ R &= 6\end{aligned}$$

First for a function ϕ we compute

$$\begin{aligned}
d\nabla^*\nabla\phi &= -\nabla_i\nabla^k\nabla_k\phi \\
&= -\nabla^k\nabla_i\nabla_k\phi - R_{ikkl}\nabla^l\phi \\
&= -\nabla^k\nabla_k\nabla_i\phi + 2\nabla_i\phi \\
&= \nabla^*\nabla d\phi + 2d\phi
\end{aligned}$$

Next

$$\begin{aligned}
\delta(\nabla^*\nabla)d\phi &= -\nabla^i\nabla^j\nabla_j\nabla_i\phi \\
&= -\nabla^i\nabla^j\nabla_i\nabla_j\phi \\
&= -\nabla^i(\nabla_i\nabla^j\nabla_j\phi + R_{jijk}\nabla_k\phi) \\
&= -(\nabla^*\nabla)^2\phi - R_{jijk}\nabla^i\nabla_k\phi \\
&= -(\nabla^*\nabla)^2\phi + 2\nabla^*\nabla\phi
\end{aligned}$$

In the fourth line we used that the curvature is parallel on S^3 . Next we compute

$$\begin{aligned}
\delta(\nabla^*\nabla)\alpha &= -\nabla^i\nabla^j\nabla_j\alpha_i \\
&= -\nabla^j\nabla_i\nabla_j\alpha_i - R_{ijjk}\nabla^k\alpha_i - R_{ijik}\nabla_j\alpha^k \\
&= -\nabla^j\nabla_i\nabla_j\alpha_i \\
&= -\nabla^j(\nabla_j\nabla_i\alpha_i + R_{ijik}\alpha^k) \\
&= \nabla^*\nabla\delta\alpha - R_{ijik}\nabla^j\alpha^k \\
&= \nabla^*\nabla\delta\alpha - 2\delta\alpha
\end{aligned}$$

and again we have used that the curvature is parallel. Using this computation we have

$$\begin{aligned}
\delta(\nabla^*\nabla)^2\alpha &= \nabla^*\nabla\delta\nabla^*\nabla\alpha - 2\delta\nabla^*\nabla\alpha \\
&= \nabla^*\nabla(\nabla^*\nabla\delta\alpha - 2\delta\alpha) - 2(\nabla^*\nabla\delta\alpha - 2\delta\alpha) \\
&= (\nabla^*\nabla)^2\delta\alpha - 4\nabla^*\nabla\delta\alpha + 4\delta\alpha
\end{aligned}$$

Also we compute

$$\begin{aligned}
\delta L_{\alpha^*}g &= \nabla^i(\nabla_i\alpha_j + \nabla_j\alpha_i) \\
&= -\nabla^*\nabla\alpha + \nabla_j\nabla^i\alpha_i + R_{ijil}\alpha^l \\
&= -\nabla^*\nabla\alpha + \nabla\delta\alpha + 2\alpha
\end{aligned}$$

Finally we compute

$$\begin{aligned}
L_{\nabla^* \nabla \tau} g &= -\nabla_i \nabla^k \nabla_k \tau_j - \nabla_j \nabla^k \nabla_k \tau_i \\
&= -\nabla^k \nabla_i \nabla_k \tau_j - R_{ikkl} \nabla^l \tau_j - R_{ikjl} \nabla^k \tau^l \\
&\quad - \nabla^k \nabla_j \nabla_k \tau_i - R_{jkk l} \nabla^l \tau_i - R_{jkil} \nabla^k \tau^l \\
&= -\nabla^k \nabla^k \nabla_i \tau_j - 2R_{ikjl} \nabla^k \tau^l + 2\nabla_i \tau_j \\
&\quad - \nabla^k \nabla^k \nabla_j \tau_i - 2R_{jkil} \nabla^k \tau^l + 2\nabla_j \tau_i \\
&= \nabla^* \nabla L_{\tau} g + 4L_{\tau} g - 2\delta \tau g
\end{aligned}$$

Now for B a symmetric two-tensor

$$\begin{aligned}
\delta \nabla^* \nabla B &= -\nabla^i \nabla^j \nabla_j B_{ik} \\
&= -\nabla^j \nabla^i \nabla_j B_{ik} - R_{ijjl} \nabla^l B_{ik} - R_{ijil} \nabla^j B_k^l - R_{ijkl} \nabla^j B_i^l \\
&= -\nabla^j \nabla^i \nabla_j B_{ik} - R_{ijkl} \nabla^j B_i^l \\
&= -\nabla^j \nabla^i \nabla_j B_{ik} - \delta B + d \operatorname{tr} B \\
&= -\nabla^j (\nabla_j \nabla^i B_{ik} + R_{ijil} B_k^l + R_{ijkl} B_i^l) - \delta B + d \operatorname{tr} B \\
&= \nabla^* \nabla \delta B - R_{ijil} \nabla^j B_k^l - R_{ijkl} \nabla^j B_{il} - \delta B + d \operatorname{tr} B \\
&= \nabla^* \nabla \delta B - 4\delta B + 2d \operatorname{tr} B
\end{aligned}$$

Using this computation we have

$$\begin{aligned}
\delta (\nabla^* \nabla)^2 &= \nabla^* \nabla \delta \nabla^* \nabla B - 4\delta \nabla^* \nabla B \\
&= \nabla^* \nabla (\nabla^* \nabla \delta B - 4\delta B) - 4(\nabla^* \nabla \delta B - 4\delta B) \\
&= (\nabla^* \nabla)^2 \delta B - 8\nabla^* \nabla \delta B + 16\delta B
\end{aligned}$$

□

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